

RECONSTRUCTION OF LAMÉ MODULI AND DENSITY AT THE BOUNDARY ENABLING DIRECTIONAL ELASTIC WAVEFIELD DECOMPOSITION*

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Abstract. We consider the inverse boundary value problem for the system of equations describing elastic waves in isotropic media on a bounded domain in \mathbb{R}^3 via a finite-time Laplace transform. The data are the dynamical Dirichlet-to-Neumann map. More precisely, using the full symbol of the transformed Dirichlet-to-Neumann map viewed as a semiclassical pseudodifferential operator, we give an explicit reconstruction of both Lamé parameters and the density, as well as their derivatives, at the boundary. We also show how this boundary reconstruction leads to a decomposition of incoming and outgoing waves.

Key words. inverse boundary value problem, layer stripping, elastic waves, isotropy

AMS subject classifications. 35R30, 35L10

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1. Introduction. We let $\Omega \subset \mathbb{R}^3$ be a bounded domain with a smooth boundary $\partial\Omega$. We consider the following initial boundary value problem for the system of equations describing elastic waves:

$$(1.1) \quad \begin{cases} \rho \partial_t^2 u = \operatorname{div}(\mathbf{C}\varepsilon(u)) =: Lu & \text{in } \Omega_T = \Omega \times (0, T), \\ u = f & \text{on } \Sigma = \partial\Omega \times (0, T), \\ u(x, 0) = \partial_t u(x, 0) = 0 & \text{in } \Omega \end{cases}$$

with $f(x, 0) = 0$ and $\frac{\partial}{\partial t} f(x, 0) = 0$ for $x \in \partial\Omega$. Here, u denotes the displacement vector and $\varepsilon(u) = (\varepsilon_{ij}(u)) = (\nabla u + (\nabla u)^T)/2$ the linear strain tensor which is the symmetric part of ∇u . Furthermore, $\mathbf{C} = \mathbf{C}(x) = (\dot{C}_{ijkl}(x))$ is the elasticity tensor and ρ is the density of mass. We assume that \mathbf{C} is isotropic, that is,

$$(1.2) \quad \dot{C}_{ijkl}(x) = \lambda(x)\delta_{ij}\delta_{kl} + \mu(x)(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})$$

with Kronecker's delta δ_{ij} and Lamé moduli $\lambda, \mu \in C^\infty(\overline{\Omega})$ such that $\mu > 0$ and $\lambda + 2\mu > 0$ on $\overline{\Omega}$. Also, $\rho \in C^\infty(\overline{\Omega})$ and $\rho > 0$ on $\overline{\Omega}$.

The hyperbolic or dynamical Dirichlet-to-Neumann map (DN map) Λ_T is defined according to

$$(1.3) \quad \Lambda_T : H^2(\Sigma) \ni f \mapsto \partial_L u := (\mathbf{C}\varepsilon(u))\nu|_{\partial\Omega} \in C([0, T], H^{1/2}(\partial\Omega)),$$

where u is the solution of (1.1), $\mathbf{C}\varepsilon(u)$ is a 3×3 matrix with its (i, k) component $(\mathbf{C}\varepsilon(u))_{ik}$ given by $(\mathbf{C}\varepsilon(u))_{ik} = \sum_{j,l=1}^3 \dot{C}_{ijkl}\varepsilon_{jl}(u)$, ν is the outward unit normal to $\partial\Omega$. Physically, $\partial_L u$ denotes the traction at $\partial\Omega$.

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In this paper, we consider the inverse problem of recovering λ, μ, ρ , as well as all their derivatives, at the boundary $\partial\Omega$ from Λ_T . Our major result for this inverse problem is as follows.

THEOREM 1.1. *The DN map Λ_T identifies λ, μ, ρ and all their derivatives on $\partial\Omega$ uniquely. There is an explicit reconstruction procedure for these identifications.*

Remark 1.2. Since the procedure we present to recover λ, μ, ρ and their derivatives at $\partial\Omega$ is local, we also have a localized version of Theorem 1.1 with partial boundary data. That is, in the definition of Λ_T we can replace $(\mathbf{C}\varepsilon(u))\nu|_{\partial\Omega}$ by $(\mathbf{C}\varepsilon(u))\nu|_{\Gamma_0}$ and confine the boundary sources, f , to those with $\text{supp } f(\cdot, t) \subset \overline{\Gamma_0}$ ($t \in (0, T)$), where Γ_0 is a relatively open subset of $\partial\Omega$. We can recover λ, μ, ρ and all their derivatives at Γ_0 . As an additional explanation for this which should be given later, see the last paragraph of section 2.

The uniqueness of the inverse problem considered here was established by Rachele [16], but giving a procedure to recover the parameters, (λ, μ, ρ) , has been left open for over 15 years. One of the complications is the occurrence of two metrics in the dynamical system of equations that cannot be straightforwardly separated at the boundary. Indeed, the usual special solutions including high-frequency asymptotic ones or progressive wave solutions based on polarization decoupling are coupled at the boundary.

Moreover, the determination of (λ, μ, ρ) on the boundary implies the following.

COROLLARY 1.3. *In addition to the conditions appearing in Theorem 1.1, let λ, μ, ρ be real analytic in the neighborhood of $\overline{\Omega}$. Then Λ_T determines uniquely λ, μ, ρ on $\overline{\Omega}$.*

A brief remark for this corollary should be given. Although there is a more general result by Rachele [16, 17], the context and argument of deriving this corollary differs from those of Rachele's.

For the static elastic inverse boundary value problem, an explicit reconstruction of λ and μ at the boundary from the full symbol of the static DN map was obtained [13, 14]. For the reconstruction of a transversely isotropic elasticity tensor, see [15]. The approach was originally developed by Sylvester and Uhlmann [21] for the electrical impedance tomography problem. The approach is also applied to Maxwell's equations [20, 12]. We generalize this type of reconstruction to dynamical elastic inverse boundary value problems. We note that our procedure is quite general and can be extended from isotropy to anisotropy with certain symmetries, which is the subject of a forthcoming paper.

The key component of the reconstruction is a connection between Λ_T and the asymptotic expansion of the DN map, a semiclassical pseudodifferential operator Λ^h say, for some elliptic system of equations containing a small parameter h via a finite-time Laplace transform. Ikehata has been using the finite Laplace transform effectively to develop his enclosure method both for parabolic equations and hyperbolic equations; see [8] and reference therein. For the convenience of our description, the partial differential operator of this system is referred to as \mathcal{M} . We will identify (λ, μ, ρ) and all of their derivatives from Λ^h by factorizing \mathcal{M} into the product of two first order semiclassical pseudodifferential operators with small parameter $h = \frac{1}{\tau}$, where τ is nothing but the Laplace variable of this transform. Also this factorization is nothing but the one used to provide the up/down going decomposition of waves which is equivalent to the incoming/downgoing decomposition of waves in the Laplace domain (see Figure 1). Further this decomposition can be linked to the corresponding de-

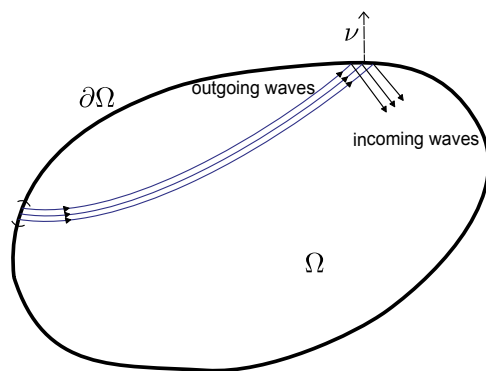


FIG. 1. Incoming/outgoing waves.

composition in the space time domain. We will briefly discuss this connection more precisely at the end of this paper. The up/down going decomposition for scalar waves is discussed in [19], and for elastic waves in [6]. In this paper, we connect the up/down going decomposition with the DN map.

Concerning our inverse problem, there are two byproducts of the factorization. One is the explicit form of the principal symbol of $\Lambda^h(s)$ and the relation of its s -derivatives to the nonprincipal symbols of $\Lambda^h(s)$. Here $\Lambda^h(s)$ is the DN map defined likewise Λ^h on the boundary $\Gamma(s)$ of the subdomain $\Omega(s) = \{x \in \Omega : \text{dist}(x, \partial\Omega) > s\}$ of Ω with $0 < s \ll 1$. The other is that a modification $\hat{\Lambda}(s)$ (cf. (3.20)) of $\Lambda^h(s)$ satisfies a Riccati-type equation. By solving the Riccati-type equation, which is an initial value problem, one can propagate the data into at least a thin layer near the boundary. This technique is known as invariant embedding in extensive geophysics literature [2, 4, 5, 6, 11].

Knowing all the derivatives of the coefficients at $\partial\Omega$ and using the Riccati equation for $\hat{\Lambda}(s)$, we can generate an approximation of $\Lambda_T(s)$ on $\Gamma(s)$. Then, we can get an approximation of λ, μ, ρ and their derivatives at $\Gamma(s)$. Repeating this process, leads to an approximation for λ, μ, ρ layer by layer in the interior of Ω , using the DN map Λ_T as the data. The associated algorithm is called layer stripping. The layer stripping was first developed for the electrical impedance tomography problem in [3, 18]. Nakamura, Tanuma and Uhlmann [15] developed such an algorithm for the static elastic inverse boundary value problem in the case of transverse isotropy.

The key application of the problem we are considering is (reflection) seismology. In actual seismic acquisition, raw vibroseis data are modeled by the local ND map: The boundary values are given by the normal traction underneath the base plate of a vibroseis and are zero (free surface) elsewhere, while the particle displacement (in fact, velocity) is measured by geophones located in a subset of the boundary (Earth's surface) (see [1]). Although, the local DN map and Neumann-to-Dirichlet (ND) map do not have the same information, the transformed ND map and transformed DN map are microlocally inverse to each other. Since we are only dealing with the symbol of the transformed DN map, the results of this paper apply to the practical setting.

The remainder of this paper is organized as follows. In section 2, we give an asymptotic identity which connects Λ_T with Λ^h via the finite-time Laplace transform introducing variable $\tau = \frac{1}{h}$. Based on this identity, we only have to find a reconstruction procedure using Λ^h . In a similar way, DN maps $\Lambda^h(s)$ are defined on each $\Gamma(s)$. In section 3, the full symbol of semiclassical pseudodifferential operator $\Lambda^h(s)$ is

analyzed using the factorization of \mathcal{M} . Also, as a byproduct of the factorization, we develop a layer stripping algorithm. Section 4 is devoted to giving a procedure and formulas for the reconstruction of (λ, μ, ρ) and all their derivatives, from the explicit form of the principal symbol of $\Lambda^h(s)$ at each $\Gamma(s)$ for $0 \leq s \ll 1$ in terms of the boundary normal coordinates associated with $\Gamma(s)$. In the final section, section 5, we will discuss another implication of the factorization. That is we give the aforementioned link between the outgoing/incoming decomposition of waves in the Laplace domain and that in the space time domain.

2. Reduction to an elliptic boundary value problem with a small parameter. First we introduce a family of symbol classes for semiclassical pseudodifferential operators. Let $A(\cdot, \cdot; \cdot) : \mathbb{R}^{2n} \times (0, h_0) \rightarrow \mathbb{C}^{\tilde{q} \times \tilde{q}}$ be a function that is smooth in $(x, \xi) \in \mathbb{R}^{2n}$ depending on $h \in (0, h_0]$ with a small $h_0 > 0$. We say that for $m \in \mathbb{R}$, A belongs to a symbol class $\mathcal{S}(m)$ if for any $\alpha, \beta \in \mathbb{Z}_+^n$, there exists a constant $C_{\alpha, \beta} > 0$ such that

$$|D_x^\alpha D_\xi^\beta A(x, \xi; h)| \leq C_{\alpha, \beta} \langle \xi \rangle^m, \quad (x, \xi) \in \mathbb{R}^{2n}, \quad h \in (0, h_0],$$

where $\mathbb{Z}_+ = \mathbb{N} \cup \{0\}$, $\langle \xi \rangle = \sqrt{1 + |\xi|^2}$.

We say that $A \in \mathcal{S}(m)$ is called a classical symbol if for any $\alpha \in \mathbb{Z}_+^n$ and $N = 1, 2, \dots$, there exists a constant $C_{\alpha, N} > 0$ such that

$$\left| \partial_x^\alpha \left(A - \sum_{j=0}^{N-1} h^j A_j \right) \right| \leq C_{\alpha, N} h^N \langle \xi \rangle^m, \quad (x, \xi) \in \mathbb{R}^{2n}$$

with each $A_j \in \mathcal{S}(m)$ independent of h , and write

$$A \sim \sum_{j=0}^{\infty} h^j A_j \quad \text{mod } \mathcal{O}(h^\infty \mathcal{S}(m)).$$

A_0 and $\sum_{j=0}^{\infty} h^j A_j$ are called the principal symbol $\sigma(A)$ and full symbol $\tilde{\sigma}(A)$ of A , respectively. In this paper we only consider classical symbols. We also denote for $A, A' \in \mathcal{S}(m)$,

$$A \sim A' \quad \text{mod } \mathcal{O}(h^k \mathcal{S}(m)),$$

if $A - A' = h^k B$ with $B \in \mathcal{S}(m)$.

For this family of symbol classes $\mathcal{S}(m)$ with $m \in \mathbb{R}$, we use the standard theory of semiclassical pseudodifferential operators (see [10, 23]). Also we use the above notations and terminologies for symbol classes associated with semiclassical pseudodifferential operators on compact manifolds. In the sequel \tilde{q} will be either $\tilde{q} = 3$ or $\tilde{q} = 6$ which can be easily noticed in context. Furthermore, we abuse the notation $\mathcal{O}(h^k \mathcal{S}(m))$ to use it also for the associated semiclassical pseudodifferential operators.

Now we consider the following boundary value problem:

$$(2.1) \quad \begin{cases} \rho v - h^2 \operatorname{div}(\mathbf{C}\varepsilon(v)) = 0 & \text{in } \Omega, \\ v = \varphi & \text{on } \partial\Omega \end{cases}$$

with a small (real-valued) parameter $h \in (0, h_0]$. We define the corresponding DN map for (2.1) according to

$$\Lambda^h : H^2(\partial\Omega) \ni \varphi \mapsto h\partial_L v = h(\mathbf{C}\varepsilon(v))\nu|_{\partial\Omega} \in H^1(\partial\Omega),$$

where v solves (2.1). In a similar fashion, we define $\Lambda^h(s)$ by replacing Ω by $\Omega(s)$ and $\partial\Omega$ by $\Gamma(s)$ while replacing φ by $\psi \in H^2(\Gamma(s))$ (cf. (3.1)) emphasizing the s dependence. Of course, $\Lambda^h = \Lambda^h(0)$. We note that Λ^h and $\Lambda^h(s)$ are semiclassical pseudodifferential operators belonging to the class $\mathcal{S}(1)$ with $n = 2$.

We show that we can obtain the full symbol of Λ^h from Λ_T via a finite-time Laplace transform. First we introduce the finite-time Laplace transform $w \in H^2(\Omega)$ of $u \in C([0, T]; H^2(\Omega))$ by

$$w(x, \tau) = (\mathcal{L}_T u)(x, \tau) = \int_0^T u(x, t) e^{-\tau t} dt \text{ with } \tau > 0.$$

In order to establish the connection between Λ_T and Λ , we let $\chi(t) = t^2$ ($t \in [0, T]$) and define

$$\tilde{\Lambda}_T : H^2(\partial\Omega) \rightarrow H^1(\Sigma)$$

by

$$\tilde{\Lambda}_T \phi = \Lambda_T(\chi\phi).$$

For any $\phi \in H^2(\partial\Omega)$, let u solve (1.1) with boundary value $f = \chi\phi$. By the estimates for solutions of hyperbolic system (1.1), we have

$$\|\partial_t^j u(\cdot, T)\|_{H^{2-j}(\Omega)} = \mathcal{O}(\|\chi\phi\|_{H^2(\Sigma)}) = \mathcal{O}(\|\phi\|_{H^2(\partial\Omega)})$$

for $j = 0, 1$ (see [22]). Because

$$\mathcal{L}_T(\partial_t^2 u) = \tau^2 \mathcal{L}_T(u) + \partial_t u(T) e^{-\tau T} + \tau u(T) e^{-\tau T} - \partial_t u(0) - \tau u(0),$$

we have that by applying \mathcal{L}_T to both sides of (1.1), and dividing by $\tau^2 = \frac{1}{h^2}$,

$$(2.2) \quad \begin{cases} \rho w - h^2 \operatorname{div}(\mathbb{C}\varepsilon(w)) = r & \text{in } \Omega, \\ w = \mathcal{L}_T(\chi\phi) & \text{on } \partial\Omega, \end{cases}$$

where $w = \mathcal{L}_T(u)$ and r has an estimate $\|r\|_{H^1(\Omega)} \leq C e^{-\kappa\tau T} \|\phi\|_{H^2(\partial\Omega)}$ with some constant C for any given κ satisfying $0 < \kappa < 1$.

Subtracting (2.1) from (2.2) with $\varphi = \mathcal{L}_T(\chi\phi)$, we find that $z = w - v$ satisfies

$$(2.3) \quad \begin{cases} \rho z - h^2 \operatorname{div}(\mathbb{C}\varepsilon(z)) = r & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega. \end{cases}$$

Hence, we have

$$\mathcal{L}_T(\partial_L u) = \partial_L w = \partial_L v + \partial_L z$$

with

$$\|h\partial_L z\|_{H^{1/2}(\partial\Omega)} \leq C\|r\|_{H^1(\Omega)} \leq C e^{-\kappa\tau T} \|\phi\|_{H^2(\partial\Omega)}$$

by standard elliptic regularity theory. We observe that

$$h\partial_L v = h\Lambda^h(\mathcal{L}_T(\chi\phi))$$

and

$$h\mathcal{L}_T(\partial_L u) = h\mathcal{L}_T\Lambda_T(\chi\phi).$$

Thus, defining $\tilde{\mathcal{L}}_T : H^2(\partial\Omega) \rightarrow H^2(\partial\Omega)$ by $\tilde{\mathcal{L}}_T(\phi) = \mathcal{L}_T(\chi\phi)$, we can rewrite the formula above as

$$h\mathcal{L}_T\tilde{\Lambda}_T = \Lambda^h\tilde{\mathcal{L}}_T + \mathcal{O}(e^{-\kappa\tau T}).$$

Here, $\mathcal{O}(e^{-\kappa\tau T})$ denotes an operator from $H^2(\partial\Omega)$ to $H^1(\partial\Omega)$ with the estimate

$$\|\mathcal{O}(e^{-\kappa\tau T})\|_{H^2(\partial\Omega)\rightarrow H^1(\partial\Omega)} \leq Ce^{-\kappa\tau T}.$$

We note that $\tilde{\mathcal{L}}_T$ is just a multiplication by $\int_0^T t^2 e^{-\tau t} dt$, hence, it is invertible and $\tilde{\mathcal{L}}_T^{-1}$ can be estimated by $\mathcal{O}(\tau^{-3})$ for $\tau \gg 1$. Therefore,

$$h\mathcal{L}_T\tilde{\Lambda}_T\tilde{\mathcal{L}}_T^{-1} \sim \Lambda^h$$

modulo an operator in $H^2(\partial\Omega) \rightarrow H^1(\partial\Omega)$ with the estimate $\mathcal{O}(h^\infty)$. Therefore, we can obtain the full symbol of Λ^h from $\mathcal{L}_T\tilde{\Lambda}_T\tilde{\mathcal{L}}_T^{-1}$, due to what we will mention in the last paragraph of this section.

We note that in the above one can choose any smooth function for χ that is consistent with the initial conditions such that $\int_0^T \chi(t)e^{-\tau t} dt$ behaves polynomially in τ .

The full symbol of a semiclassical pseudodifferential operator can be evaluated by applying to locally supported rapidly oscillating functions [23]. We note here that the analysis can be local, and thus we have Remark 1.2.

3. Analysis of the symbol of $\Lambda^h(s)$. Given a boundary point $p_0 \in \Gamma(s)$, for any $x \in \Omega(s)$ near p_0 , we use the boundary normal coordinates $x = (x^1(p), x^2(p), x^3) = (y^1, y^2, x^3) = (y', x^3)$, where $p \in \Gamma(s)$ is the nearest point to x , $x^3 = \text{dist}(x, p)$, and $(x^1(p), x^2(p))$ are the local coordinates of $\Gamma(s)$ near p_0 . Then $\Gamma(s)$ is locally given as $x^3 = s$. Let (ξ_1, ξ_2, ξ_3) , (η_1, η_2, η_3) be conormal vectors with respect to the coordinates (x_1, x_2, x_3) , (x^1, x^2, x^3) such that $\sum_{j=1}^3 \xi_j dx_j = \sum_{j=1}^3 \eta_j dx^j$. We will use the notation $\eta = (\eta_1, \eta_2, \eta_3) = (\eta', \eta_3)$. In boundary normal coordinates, (2.1) attains the form

$$(3.1) \quad \begin{cases} (\mathcal{M}v)^i = \rho g^{ik} v_k - h^2 \sum_{j,k,l=1}^3 \nabla_j (C^{ijkl} \varepsilon_{kl}(v)) = 0 \text{ in } \{x^3 > s\} \text{ for } 1 \leq i \leq 3, \\ v^i|_{x^3=s} = \psi^i, \quad 1 \leq i \leq 3, \end{cases}$$

where ∇_j is the covariant derivative with respect to $\frac{\partial}{\partial x^j}$ and $\varepsilon_{kl}(v) = 2^{-1}(\nabla_l v_k + \nabla_k v_l)$ is the linear strain tensor,

$$C^{ijkl}(x) = \sum_{a,b,c,d=1}^3 \frac{\partial x^i}{\partial x_a} \frac{\partial x^j}{\partial x_b} \frac{\partial x^k}{\partial x_c} \frac{\partial x^l}{\partial x_d} \dot{C}_{abcd}(x)$$

with $\dot{C}_{abcd}(x)$ given by (1.2). The induced metric $G(x) = (g^{ai}(x))$ is given by

$$g^{ai}(x) = \sum_{r=1}^3 \frac{\partial x^a}{\partial x_r}(x) \frac{\partial x^i}{\partial x_r}(x).$$

In terms of Jacobi matrix $J = (\partial x^a / \partial x_r; 1 \leq a, r \leq 3)$, G takes the form $G = JJ^T$.

The expression for $\Lambda^h(s)$ in boundary normal coordinates is

$$(3.2) \quad (\Lambda^h(s)\psi)^i = -h \sum_{k,l=1}^3 C^{i3kl} \varepsilon_{kl}(v), \quad 1 \leq i \leq 3,$$

at $x^3 = s$. The full symbol of $\Lambda^h(s)$ can be expanded as

$$\tilde{\sigma}(\Lambda^h(s))(y', \eta') \sim \sum_{j \leq 0} h^{-j} \lambda_{-j}(s)(y', \eta') \mod \mathcal{O}(h^\infty \mathcal{S}(1)),$$

where each $\lambda_j(s)(y', \eta') \in \mathcal{S}(1)$.

Now, we define

$$(3.3) \quad \begin{aligned} Q(x, \eta') &= \left(\sum_{j,l=1}^2 C^{ijkl}(x) \eta_j \eta_l; 1 \leq i, k \leq 3 \right), \\ R(x, \eta') &= \left(\sum_{j=1}^2 C^{ijk3}(x) \eta_j; 1 \leq i, k \leq 3 \right), \\ D(x) &= (C^{i3k3}(x); 1 \leq i, k \leq 3). \end{aligned}$$

The principal symbol, $M(x, \eta)$, of \mathcal{M} is then given by

$$(3.4) \quad M(x, \eta) = D(x) \eta_3^2 + (R(x, \eta') + R^T(x, \eta')) \eta_3 + Q(x, \eta') + \rho(x) G(x).$$

By the assumption, $M(x, \eta)$ is a positive definite matrix for $x \in \overline{\Omega}$, $\eta \in \mathbb{R}^3 \setminus 0$. Hence, for fixed (x, η') , $\det D^{-1/2} M(x, \eta) D^{-1/2} = 0$ in η_3 admits 3 roots $\eta_3 = \zeta_j$ ($j = 1, 2, 3$) with positive imaginary parts and 3 roots $\overline{\zeta_j}$ ($j = 1, 2, 3$) with negative imaginary parts. Thus, we have the following lemma.

LEMMA 3.1 (see [7]). *There is a unique factorization*

$$\tilde{M}(x, \eta) = D(x)^{-1/2} M(x, \eta) D(x)^{-1/2} = (\eta_3 - \tilde{S}_0^*(x, \eta')) (\eta_3 - \tilde{S}_0(x, \eta'))$$

with $\text{Spec}(\tilde{S}_0(x, \eta')) \subset \mathbb{C}_+$, where $\text{Spec}(\tilde{S}_0(x, \eta'))$ is the spectrum of $\tilde{S}_0(x, \eta')$. In the above,

$$\tilde{S}_0(x, \eta') := \left(\oint_{\gamma} \zeta \tilde{M}(x, \eta', \zeta)^{-1} d\zeta \right) \left(\oint_{\gamma} \tilde{M}(x, \eta', \zeta)^{-1} d\zeta \right)^{-1},$$

where $\gamma \subset \mathbb{C}_+ := \{\zeta \in \mathbb{C} : \Im \zeta := \text{imaginary part of } \zeta > 0\}$ is a continuous curve enclosing all the ζ_j ($j = 1, 2, 3$).

Then we have the following factorization of $M(x, \eta)$:

$$(3.5) \quad M(x, \eta) = (\eta_3 - S_0^*(x, \eta')) D(x) (\eta_3 - S_0(x, \eta')),$$

where

$$S_0(x, \eta') = D^{-1/2}(x) \tilde{S}_0(x, \eta') D^{1/2}(x).$$

We arrive at the following.

LEMMA 3.2. *The operator \mathcal{M} admits a factorization*

$$(3.6) \quad \begin{aligned} \mathcal{M} &= (hD_s - S^*(x, hD_{y'}; h) \\ &\quad + hK(x, hD_{y'})) D(x) (hD_s - S(x, hD_{y'}; h)), \end{aligned}$$

where $S(x, \eta'; h) \in \mathcal{S}(1)$, $K(x, \eta'; h) \in \mathcal{S}(0)$. Moreover, $hD_{y'} = (hD_{y^1}, hD_{y^2})$, $hD_{y^j} = -i h \partial / \partial y^j$ ($j = 1, 2$), and the principal symbol, $S_0(x, \eta')$, of S satisfies

$$(3.7) \quad \text{Spec}(S_0(x, \eta')) \subset \mathbb{C}_+.$$

Proof. Following (3.4), we write the full symbol $\tilde{\sigma}(\mathcal{M})$ of \mathcal{M} in the form

$$(3.8) \quad \begin{aligned} \tilde{\sigma}(\mathcal{M}) = & D(x)\eta_3^2 + (R(x, \eta') + R^T(x, \eta'))\eta_3 \\ & + Q(x, \eta') + \rho(x)G(x) + hF_0(x)\eta_3 + hF_1(x, \eta'), \end{aligned}$$

where $F_1(x, \eta') \in \mathcal{S}(1)$, and $F_0(x)$ is a matrix multiplication. We expand

$$(hD_s - S^*(x, hD_{y'}; h) + hK(x, hD_{y'}; h))D(x)(hD_s - S(x, hD_{y'}; h)),$$

yielding

$$(3.9) \quad \begin{aligned} & h(D_s D)(hD_s) - S^* D(hD_s) + hK D(hD_s) - h(D_s D)S + S^* D S \\ & - hK D S - hD(D_s S) - D S(hD_s) + D(hD_s)^2. \end{aligned}$$

Comparing (3.8) and (3.9), we find that S and K should satisfy

$$(3.10) \quad \begin{aligned} & -S^*(x, hD_{y'}; h)D(x) + hK(x, hD_{y'}; h)D(x) - D(x)S(x, hD_{y'}; h) + hD_s D(x) \\ & = R(x, hD_{y'}) + R^T(x, hD_{y'}) + hF_0(x) \end{aligned}$$

and

$$(3.11) \quad \begin{aligned} & -h(D_s D(x))S(x, hD_{y'}; h) + S^*(x, hD_{y'}; h)D(x)S(x, hD_{y'}; h) \\ & - hK(x, hD_{y'}; h)D(x)S(x, hD_{y'}; h) - hD(x)(D_s S(x, hD_{y'}; h)) \\ & = hF_1(x, hD_{y'}) + Q(x, hD_{y'}) + \rho G(x). \end{aligned}$$

Eliminating K in (3.11) by using (3.10), we get

$$(3.12) \quad (hD_s)S + S^2 + D^{-1}(R + R^T + hF_0)S + hD^{-1}F_1 + D^{-1}Q + D^{-1}\rho G = 0.$$

By the composition formula for symbols of pseudodifferential operators, we have

$$(3.13) \quad \begin{aligned} & \sum_{\alpha \geq 0} \frac{i^{|\alpha|}}{\alpha!} h^{|\alpha|} D_{\eta'}^{\alpha} S(x, \eta'; h) D_{y'}^{\alpha} S(x, \eta'; h) \\ & + \sum_{\alpha \geq 0} \frac{i^{|\alpha|}}{\alpha!} h^{|\alpha|} D_{\eta'}^{\alpha} (D^{-1}(x)(R(x, \eta') + R^T(x, \eta'))) D_{y'}^{\alpha} S(x, \eta'; h) \\ & + hD^{-1}(x)F_0(x)S(x, \eta'; h) + hD^{-1}(x)F_1(x, \eta') + D^{-1}(x)Q(x, \eta') \\ & + D^{-1}(x)\rho G + hD_s S(x, \eta'; h) = 0. \end{aligned}$$

We introduce the expansions

$$\begin{aligned} S(x, \eta'; h) & \sim \sum_{j \leq 0} h^{-j} S_j(x, \eta'), \\ K(x, \eta'; h) & \sim \sum_{j \leq 0} h^{-j} K_j(x, \eta') \end{aligned}$$

with $S_j \in \mathcal{S}(1)$, $K_j \in \mathcal{S}(0)$ for every j . We construct S via arranging terms of the same degree of h in (3.13). The terms of order $\mathcal{O}(h^0)$ give

$$(3.14) \quad \begin{aligned} & D^{-1}(x)(R(x, \eta') + R^T(x, \eta'))S_0(x, \eta') + D^{-1}(x)Q(x, \eta') \\ & + D^{-1}(x)\rho(x)G + S_0^2(x, \eta') = 0. \end{aligned}$$

Indeed, S_0 (cf. (3.5)) satisfies this equation. The terms of order $\mathcal{O}(h)$ give

$$(3.15) \quad S_0 S_{-1} + S_{-1} S_0 + D^{-1}(R + R^T)S_{-1} + \sum_{|\alpha|=1} iD_{\eta'}^{\alpha} S_0 D_{y'}^{\alpha} S_0 \\ + \sum_{|\alpha|=1} iD_{\eta'}^{\alpha} (D^{-1}(R + R^T))D_{y'}^{\alpha} S_0 + D^{-1}F_0 S_0 + D^{-1}F_1 + D_s S_0 = 0.$$

The terms which are of homogeneity of order $\mathcal{O}(h^{-j})$ for $j \leq -2$ yield

$$(3.16) \quad S_0 S_j + S_j S_0 + D^{-1}(R + R^T)S_j + \sum_{\substack{l+m=j+|\alpha| \\ |\alpha| \geq 1}} \frac{i^{|\alpha|}}{\alpha!} D_{\eta'}^{\alpha} S_l D_{y'}^{\alpha} S_m + \sum_{\substack{l+m=j \\ m, l \leq 0}} S_l S_m \\ + \sum_{|\alpha|=1} iD_{\eta'}^{\alpha} (D^{-1}(R + R^T))D_{y'}^{\alpha} S_{j+1} + D^{-1}F_0 S_{j+1} + D_s S_{j+1} = 0.$$

To confirm that (3.15) and (3.16) have solutions, we note that

$$S_j S_0 + S_0 S_j + D^{-1}(R + R^T)S_j = -D^{-1}(Q + \rho G)S_0^{-1}S_j + S_j S_0,$$

using (3.14). Since

$$\text{Spec}(S_0) \subset \mathbb{C}_+, \quad \text{Spec}(-D^{-1}(Q + \rho G)S_0^{-1}) \subset \mathbb{C}_+,$$

we can indeed solve for S_j ($j \leq -1$) in (3.15) and (3.16).

After constructing the full symbol of S , we determine the full symbol of K from (3.10). \square

PROPOSITION 3.3. *Let $\lambda_0(s)(y, \eta')$ be the principal symbol of $\Lambda^h(s)$. Then*

$$(3.17) \quad \lambda_0(s)(y', \eta') = -i(D(x)S_0(x, \eta') + R^T(x, \eta'))|_{x^3=s}.$$

Proof. For a given s ($0 < s \ll 1$), let \mathcal{T} be given such that $0 < \mathcal{T} - s \ll 1$. The parametrix $U = U(y', x^3; h)$ to the boundary value problem (2.1) satisfies locally

$$\mathcal{M}U \sim 0 \mod \mathcal{O}(h^\infty \mathcal{S}(2)) \text{ in } \mathbb{R}^2 \times [s, \mathcal{T}], \\ U|_{x^3=\mathcal{T}} = I.$$

Equation (3.7) implies that the solution operator of the factor

$$(hD_s) - S^*(x, hD_{y'}; h) + hK(x, hD_{y'}; h)$$

in the factorization (3.6), for decreasing s , is decaying of order $\mathcal{O}(h^\infty)$. Hence, U satisfies

$$((hD_s) - S(x, hD_{y'}; h))U \sim 0 \mod \mathcal{O}(h^\infty \mathcal{S}(1)).$$

Thus,

$$(hD_s)U|_{x^3=s} \sim S(x, hD_{y'}; h)U|_{x^3=s} \mod \mathcal{O}(h^\infty \mathcal{S}(1)).$$

Therefore,

$$(3.18) \quad \Lambda^h(s)(y', hD_{y'}; h) \sim -i(D(x)S(x, hD_{y'}; h) + R^T(x, hD_{y'}))|_{x^3=s} \mod \mathcal{O}(h^\infty \mathcal{S}(1)).$$

Then formula (3.17) follows immediately. \square

Next we establish the relation between the principal symbol $\lambda_0(s)$ and the lower order ones $\lambda_j(s)$, $j \leq -1$. Below, we use the notation $\text{mod } (T_s^k, h\mathcal{S}(1))$ to indicate ignoring terms in $h\mathcal{S}(1)$ and in $T_s^k = \{\text{symbol } p_s(y', \eta') \text{ which depends only on the } s\text{-derivatives of } \lambda(y', s), \mu(y', s), \rho(y', s) \text{ up to order } k\}$.

PROPOSITION 3.4. *There is a bijective linear map $W(y', s, \eta')$ on the set of 3×3 matrices which depends only on λ, μ, ρ , but not on their normal derivatives, such that*

$$(3.19) \quad \lambda_j(s)(y', \eta') \sim W(\cdot, s, \cdot)(D_s \lambda_{j+1}(s))(y', \eta') \quad \text{mod } (T_s^{-(1+j)}, h\mathcal{S}(1))$$

for any $j \leq -1$.

Proof. Throughout the proof, we read $x^3 = s$. First, we note that for $j \leq -1$

$$\lambda_j(s)(y', \eta') = -iD(x)S_j(x, \eta').$$

From (3.15), we obtain

$$S_0 S_{-1} + S_{-1} S_0 + D^{-1}(R + R^T)S_{-1} = -D_s S_0 - D^{-1}F_0 S_0 - D^{-1}F_1 \quad \text{mod } (T_s^0, h\mathcal{S}(1)).$$

Moreover,

$$F_0 = D_s D \quad \text{mod } (T_s^0, h\mathcal{S}(1))$$

and

$$F_1 = D_s R^T \quad \text{mod } (T_s^0, h\mathcal{S}(1)).$$

Hence,

$$\begin{aligned} S_0 S_{-1} + S_{-1} S_0 + D^{-1}(R + R^T)S_{-1} &= -D_s S_0 - D^{-1}(D_s D)S_0 - D^{-1}(D_s R^T) \\ &= -D^{-1}D_s(DS_1 + R^T) \\ &= iD^{-1}D_s \lambda_1(s) \quad \text{mod } (T_s^0, h\mathcal{S}(1)). \end{aligned}$$

Following the proof of Lemma 3.2, we find that $\lambda_{-1}(s)$ satisfies

$$(Q + \rho G)S_0^{-1}D^{-1}\lambda_{-1}(s) - \lambda_{-1}(s)S_0 = D_s \lambda_0(s) \quad \text{mod } (T_s^0, h\mathcal{S}(1)).$$

We note that

$$\text{Spec}(S_0) \subset \mathbb{C}_+, \quad \text{Spec}(-(Q + \rho G)S_0^{-1}D^{-1}) \subset \mathbb{C}_+.$$

Hence, defining $W(x, \eta')(Y)$ as the solution X of

$$(Q + \rho G)S_1^{-1}D^{-1}X - XS_0 = Y,$$

we obtain

$$\lambda_{-1}(s)(y', \eta') = W(\cdot, s, \cdot)D_s \lambda_0(s)(y', \eta') \quad \text{mod } (T_s^0, h\mathcal{S}(1)).$$

For $j \leq -2$, S_j contains s -derivatives of λ, μ, ρ up to order $-j$. Then, inductively, we get

$$(Q + \rho G)S_0^{-1}D^{-1}\lambda_j(s) - \lambda_j(s)S_0 = D_s \lambda_{j+1}(s) \quad \text{mod } (T_s^{-(1+j)}, h\mathcal{S}(1)).$$

Thus, we have proved the claim. \square

We conclude this section by presenting the Riccati equation that $\Lambda(s)$ satisfies.

COROLLARY 3.5. *Define*

$$(3.20) \quad \hat{\Lambda}(s) = iD^{-1}\Lambda^h(s);$$

$\hat{\Lambda}(s)$ satisfies, mod $\mathcal{O}(h^\infty \mathcal{S}(1))$, the Riccati equation

$$(3.21) \quad hD_s \hat{\Lambda}(s) + J_1(s) \hat{\Lambda}(s) + \hat{\Lambda}(s) K_1(s) + \hat{\Lambda}(s)^2 + F_2(s) = 0 \quad (0 \leq s \ll 1),$$

where

$$J_1(s) = D^{-1}(R + hF_0), \quad K_1(s) = -D^{-1}R^T,$$

and

$$F_2(s) = -hD_s(D^{-1}R^T) - D^{-1}(R + R^T + hF_0)D^{-1}R^T + D^{-1}(hF_1 + Q + \rho G) + (D^{-1}R^T)^2,$$

with $x^3 = s$.

Proof. This follows straightforwardly from (3.12) and (3.18). \square

Invoking a forward Euler scheme to solve the Riccati equation, we obtain an approximate propagation of the boundary data into the interior of Ω , layer by layer. With the explicit reconstruction that will be presented in the next section, we obtain formally a layer-stripping algorithm for our inverse boundary value problem. The Riccati-type equation is expected to be highly unstable, especially for high-frequency modes [18]. So the propagation of the DN map will deteriorate. This reveals the ill-posedness of the problem of recovering the parameter in the interior.

4. Reconstruction of the Lamé parameters and density. We present the reconstruction of (λ, μ, ρ) , as well as all their derivatives at $x^3 = s$ from the full symbol of $\Lambda(s)$. We first consider the principal symbol of operator $-h^2 \sum_{j,k,l=1}^3 \nabla_j (C^{ijkl} \varepsilon_{kl}(v))$. By the transformation rule of tensors, we have

$$(4.1) \quad N := \left(\sum_{j,l=1}^3 C^{ijkl} \eta_j \eta_l; 1 \leq i, k \leq 3 \right) = J \dot{N} J^T$$

with $\dot{N} = \left(\sum_{j,l=1}^3 \dot{C}_{ijkl} \xi_j \xi_l; 1 \leq i, k \leq 3 \right).$

For any x near $\Gamma(s)$, we choose a unit vector $n(x) = (n_1, n_2, n_3) \in \mathbb{R}^3$ depending smoothly on x . Then any $\xi = (\xi_1, \xi_2, \xi_3) \in \mathbb{R}^3$ can be written as $\xi = qn(x) + m(x, \xi)$ for some $q \in \mathbb{R}$ and $(m_1, m_2, m_3) =: m \perp n$. We define $\dot{D} = \dot{D}(x)$, $\dot{R} = \dot{R}(x, \xi)$, $\dot{Q} = \dot{Q}(x, \xi)$ as follows:

$$(4.2) \quad \begin{aligned} \dot{D} &= \left(\sum_{j,l=1}^3 \dot{C}_{ijkl} n_j n_l; 1 \leq i, k \leq 3 \right), \\ \dot{R} &= \left(\sum_{j,l=1}^3 \dot{C}_{ijkl} m_j n_l; 1 \leq i, k \leq 3 \right), \\ \dot{Q} &= \left(\sum_{j,l=1}^3 \dot{C}_{ijkl} m_j m_l; 1 \leq i, k \leq 3 \right). \end{aligned}$$

We have (compare with (3.4))

$$(4.3) \quad \dot{M} = \dot{D}q^2 + \left(\dot{R} + (\dot{R})^T \right) q + \dot{Q} + \rho,$$

and a smooth factorization according to (3.5). More precisely, there exists a unique $\dot{S}_0 = \dot{S}_0(x, \xi)$ depending smoothly on $x \in \overline{\Omega}$, and homogeneous of degree one with respect to ξ , such that

$$(4.4) \quad \dot{M} = \left(q - (\dot{S}_0)^* \right) \dot{D}(q - \dot{S}_0), \quad \text{Spec}(\dot{S}_0) \subset \mathbb{C}_+.$$

We let the direction of $n(x)$ be aligned with the x^3 axis. Using (4.1) we find that

$$(4.5) \quad M = \left(q - J(\dot{S}_0)^* J^{-1} \right) \left(J \dot{D} J^T \right) \left(q - (J^T)^{-1} \dot{S}_0 J^T \right).$$

Since the linear mapping defined by the matrix $(J^T)^{-1}$ preserves the orthogonality $m(x, \xi) \perp n(x)$ and the length of $n(x)$, we have $q = \eta_3$. We also have $D = J \dot{D} J^T$. Hence, by the uniqueness of factorization (3.5),

$$(4.6) \quad S_0 = (J^T)^{-1} \dot{S}_0 J^T.$$

Combining this with (3.17) and the tensorial transformation $R = J \dot{R} J^T$ of \dot{R} , we obtain

$$(4.7) \quad \lambda_0(s) = -iJ \left(\dot{D} \dot{S}_0 + \dot{R}^T \right) J^T.$$

Due to the isotropy of the elasticity tensor, a rotation of coordinates (x_1, x_2, x_3) does not affect the form \dot{D} , \dot{S}_1 , \dot{R} and \dot{Q} . Hence, for any x , we just assume $\xi = (\xi', \xi_3)$ and $n(x) = (0, 0, 1)$. Then $m(x, \xi) = \xi'$, so $\dot{S}_0(x, \xi)$ depends only on (x, ξ') ,

$$\dot{S}_0(x, \xi) = \dot{A}(x, \xi') + i\dot{B}(x, \xi')$$

and

$$\dot{D} = \text{diag}(\mu, \mu, \lambda + 2\mu), \quad \dot{A} = \dot{D}^{-1/2} \tilde{A} \dot{D}^{1/2}, \quad \dot{B} = \dot{D}^{-1/2} \tilde{B} \dot{D}^{1/2}$$

in which

$$(4.8) \quad \tilde{A} = P \begin{pmatrix} 0 & 0 & -\alpha_1 \\ 0 & 0 & 0 \\ -\alpha_2 & 0 & 0 \end{pmatrix} P^*, \quad \tilde{B} = P \begin{pmatrix} a & 0 & 0 \\ 0 & b & 0 \\ 0 & 0 & c \end{pmatrix} P^*$$

with

$$P = P(\xi') = \begin{pmatrix} \xi_1 |\xi'|^{-1} & \xi_2 |\xi'|^{-1} & 0 \\ \xi_2 |\xi'|^{-1} & -\xi_1 |\xi'|^{-1} & 0 \\ 0 & 0 & 1 \end{pmatrix},$$

$$\alpha_1 = \frac{(\lambda + \mu) |\xi'|}{\sqrt{\mu(\lambda + 2\mu)}} \frac{1}{1 + \gamma}, \quad \alpha_2 = \gamma \alpha_1, \quad b = \sqrt{\frac{\mu |\xi'|^2 + \rho}{\mu}},$$

$$c = \frac{1}{1 + \gamma} \sqrt{(1 + \gamma)^2 \frac{\mu |\xi'|^2 + \rho}{\lambda + 2\mu} - \frac{(\lambda + \mu)^2 |\xi'|^2}{\mu(\lambda + 2\mu)}}, \quad a = \gamma c,$$

and

$$\gamma = \sqrt{\frac{((\lambda + 2\mu)|\xi'|^2 + \rho)(\lambda + 2\mu)}{\mu(\mu|\xi'|^2 + \rho)}}.$$

We substitute $\xi' = [|\xi'|, 0]^T$ and identify ξ' with $|\xi'|$; then, after some calculations, we find that

$$\begin{aligned} \lambda_0(s)(y, \eta) \\ = -iJ \left(\dot{D}(x) \dot{S}_0(x, \xi') + \dot{R}^T(x, \xi') \right) J^T = J \dot{\Lambda}_0 J^T \end{aligned}$$

with

$$\begin{aligned} \dot{\Lambda}_1 &= (\dot{\lambda}_1^{(ik)}; 1 \leq i, k \leq 3) \\ &= \begin{pmatrix} a\mu & 0 & i\alpha_1 \sqrt{\mu(\lambda + 2\mu)} - i\mu|\xi'| \\ 0 & b\mu & 0 \\ i\alpha_2 \sqrt{\mu(\lambda + 2\mu)} - i\lambda|\xi'| & 0 & c(\lambda + 2\mu) \end{pmatrix}. \end{aligned}$$

Since J is known and independent of λ , μ , ρ , we only need to consider $\dot{\Lambda}_1$ for recovering λ , μ , ρ and their derivatives at $x^3 = s$.

First step. We will recover λ, μ, ρ . Note that $\dot{\lambda}_0^{(22)}(x, \xi') = \sqrt{|\xi'|^2 \mu^2 + \rho \mu}$. Then we can first get μ and ρ as follows. Observe that

$$\dot{\lambda}_0^{(22)}(x, \sqrt{2}c_0^{-1})^2 - \dot{\lambda}_0^{(22)}(x, c_0^{-1})^2 = c_0^{-2} \mu^2$$

for any scaling constant $c_0 > 0$. Hence we set $c_0 = 1$ in the rest of this section. Then we find

$$\mu = \sqrt{\dot{\lambda}_0^{(22)}(x, \sqrt{2}c_0^{-1})^2 - \dot{\lambda}_1^{(22)}(x, c_0^{-1})^2}$$

and

$$\rho = \frac{1}{\mu} \left(\dot{\lambda}_0^{(22)}(x, c_0^{-1})^2 - \mu^2 \right).$$

For λ , first notice that we have

$$\frac{\dot{\lambda}_0^{(11)}(x, c_0^{-1})^2}{\dot{\lambda}_0^{(33)}(x, c_0^{-1})^2} = \frac{(\lambda + 2\mu + \rho)\mu}{(\mu + \rho)(\lambda + 2\mu)}.$$

Since we have already computed μ and ρ , we get

$$\frac{\lambda + 2\mu + \rho}{\lambda + 2\mu} = 1 + \frac{\rho}{\lambda + 2\mu},$$

and then obtain λ .

Second step. We recover $\partial\lambda$, $\partial\mu$, $\partial\rho$ of λ , μ , ρ ; we first note that

$$(4.9) \quad 2\mu\partial\mu = \partial \left(\dot{\lambda}_0^{(22)}(x, \sqrt{2}c_0^{-1})^2 - \dot{\lambda}_1^{(22)}(x, c_0^{-1})^2 \right)$$

from which we can recover $\partial\mu$. Then from

$$(4.10) \quad \rho\partial\mu + \mu\partial\rho = \partial \left(\dot{\lambda}_1^{(22)}(x, c_0^{-1})^2 - \mu^2 \right)$$

we recover $\partial\rho$. Finally we recover $\partial\lambda$ from

$$(4.11) \quad \partial\lambda + 2\partial\mu = \partial \left(\left(\frac{\mu + \rho}{\mu} \frac{\dot{\lambda}_1^{(11)}(x, c_0^{-1})^2}{\dot{\lambda}_1^{(33)}(x, c_0^{-1})^2} - 1 \right)^{-1} \rho \right).$$

Final step. We recover higher order derivatives of λ, μ, ρ . Differentiating equations (4.9)–(4.11) $k-1$ times, we obtain linear equations for $\partial^k\mu, \partial^k\lambda$, and $\partial^k\rho$. The coefficients for them are the same as those for $\partial\lambda, \partial\mu, \partial\rho$ in (4.9)–(4.11). Thus we can recover $\partial^k\mu, \partial^k\lambda, \partial^k\rho$, using $\partial^j\dot{\Lambda}_0(s)(y, \eta)$ for $j = 1, 2, \dots, k$, and $\partial^j\mu, \partial^j\lambda, \partial^j\rho$ for $j = 1, 2, \dots, k-1$. So we can recover all the derivatives of λ, μ, ρ recursively.

5. Further implications. In this section, we show how to get a decomposition into incoming/outgoing waves via the factorization (3.5). In this section, we take $T = \infty$ and allow τ to take complex values. We assume τ is in the set

$$\Pi_0 = \{\tau \in \mathbb{C}; \Re\tau := \text{real part of } \tau > 0\}.$$

For a bounded $u \in C(0, \infty; H^2(\Omega))$, we introduce the Laplace transform \mathcal{L} for $\tau \in \Pi_0$:

$$(\mathcal{L}u)(x, \tau) = \int_0^\infty e^{-\tau t} u(x, t) dt,$$

where $(\mathcal{L}u)(\cdot, \tau) \in H^2(\Omega)$. The inverse Laplace transform is given by

$$(\mathcal{L}^{-1}v)(x, t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} e^{\tau t} v(x, \tau) d\tau$$

with $\gamma \in \Pi_0$.

Applying the above Laplace transform to (1.1), we get

$$\mathcal{M}v = \rho \hat{\tau}^2 g^{ik} v_k - h^2 \sum_{j,k,l=1}^3 \nabla_j (C^{ijkl} \varepsilon_{kl}(v)) = 0$$

with $h = \frac{1}{|\tau|}$, $\hat{\tau} = \frac{\tau}{|\tau|}$. \mathcal{M} can be viewed as a semiclassical pseudodifferential operator with a small parameter $h = \frac{1}{|\tau|}$.

In a neighborhood of the boundary, we rewrite the above equation up to the leading order terms in the following form,

$$(5.1) \quad \begin{aligned} & hD_s \begin{pmatrix} v \\ hD_s v \end{pmatrix} \\ & \sim \begin{pmatrix} 0 & 1 \\ -D^{-1}(Q + \rho G\tau^2) & -D^{-1}(R + R^T) \end{pmatrix} \begin{pmatrix} v \\ hD_s v \end{pmatrix} \mod \mathcal{O}(hS(2)). \end{aligned}$$

Denote $M(x, \hat{\tau}, \eta)$ to be the principal symbol of \mathcal{M} ; as in (3.5). In a neighborhood of the boundary, we have the factorization

$$(5.2) \quad M(x, \hat{\tau}, \eta) = (\eta_3 - S_0^-(x, \hat{\tau}, \eta')) D(x) (\eta_3 - S_0^+(x, \hat{\tau}, \eta'))$$

for $\tau \in \Pi_0$. Here, similarly to (4.8),

$$(5.3) \quad S_0^+(x, \hat{\tau}, \eta') = (J^T)^{-1} \dot{D}^{-1/2} (\tilde{A} + i\tilde{B}) \dot{D}^{1/2} J^T$$

and

$$(5.4) \quad S_0^-(x, \hat{\tau}, \eta') = J \dot{D}^{1/2} \left(\tilde{A}^T - i \tilde{B}^T \right) \dot{D}^{-1/2} J^{-1},$$

where

$$(5.5) \quad \tilde{A}(x, \hat{\tau}, \eta') = P \begin{pmatrix} 0 & 0 & -\alpha_1 \\ 0 & 0 & 0 \\ -\alpha_2 & 0 & 0 \end{pmatrix} P^*, \quad \tilde{B}(x, \hat{\tau}, \eta') = P \begin{pmatrix} a & 0 & 0 \\ 0 & b & 0 \\ 0 & 0 & c \end{pmatrix} P^*$$

with

$$\alpha_1 = \frac{(\lambda + \mu)|\xi'|}{\sqrt{\mu(\lambda + 2\mu)}} \frac{1}{1 + \gamma}, \quad \alpha_2 = \gamma \alpha_1, \quad b = \sqrt{\frac{\mu|\xi'|^2 + \rho \hat{\tau}^2}{\mu}},$$

$$c = \frac{1}{1 + \gamma} \sqrt{(1 + \gamma)^2 \frac{\mu|\xi'|^2 + \rho \hat{\tau}^2}{\lambda + 2\mu} - \frac{(\lambda + \mu)^2 |\xi'|^2}{\mu(\lambda + 2\mu)}}, \quad a = \gamma c,$$

and

$$\gamma = \sqrt{\frac{((\lambda + 2\mu)|\xi'|^2 + \rho \hat{\tau}^2)(\lambda + 2\mu)}{\mu(\mu|\xi'|^2 + \rho \hat{\tau}^2)}}.$$

In all formulas for $a, b, c, \alpha_1, \alpha_2, \gamma$, \sqrt{z} is defined on $\mathbb{C} \setminus (-\infty, 0]$ with $\Re \sqrt{z} > 0$. Indeed α_1, γ , and b are well-defined, and $\Re \gamma > 0$. Furthermore, we note that if $\Im \tau > 0$, then $\Im \left\{ \frac{\mu|\xi'|^2 + \rho \hat{\tau}^2}{\lambda + 2\mu} \right\} > 0$ and $\Im \gamma < 0$, while $\Im(1 + \gamma)^2 < 0$. Thus c is well-defined for $\Im \tau > 0$. Similarly, we can verify c is well-defined for $\Im \tau < 0$.

We now show that

$$(5.6) \quad \text{Spec}(S_0^+) \subset \mathbb{C}_+, \quad \text{Spec}(S_0^-) \subset \mathbb{C}_-$$

for any $\tau \in \Pi_0$. The spectrum of S_0^+ unified with the spectrum of S_0^- are the roots of $\det(M(x, \hat{\tau}, \eta))$ as a polynomial in η_3 for $\tau \in \Pi_0$. To obtain statement (5.6), we only need to show that there are no real roots of $\det(M(x, \hat{\tau}, \eta))$ for any $\tau \in \Pi_0$. Then because (5.6) holds true for positive τ , the eigenvalues of S_0^+ or S_0^- cannot intersect the real line. Roots of $\det(M(x, \hat{\tau}, \eta))$, which are same as the roots of $\det(\dot{M})$ in q , satisfy

$$\mu q^2 + \mu|\xi'|^2 + \rho \hat{\tau}^2 = 0$$

or

$$((\lambda + 2\mu)q^2 + \mu|\xi'|^2 + \rho \hat{\tau}^2) (\mu q^2 + (\lambda + 2\mu)|\xi'|^2 + \rho \hat{\tau}^2) - (\lambda + \mu)^2 q^2 |\xi'|^2 = 0.$$

For $\tau \in (0, \infty)$, we have already concluded that q could not be real. For nonreal τ , if there is a real q satisfying the above equations, then

$$\Im \hat{\tau}^2 = 0,$$

which is not possible.

Let

$$P = \begin{pmatrix} (S_0^- - S_0^+)^{-1} S_0^- & -(S_0^- - S_0^+)^{-1} \\ -(S_0^- - S_0^+)^{-1} S_0^+ & (S_0^- - S_0^+)^{-1} \end{pmatrix}$$

and denote its inverse by P^{-1} which is given by

$$P^{-1} = \begin{pmatrix} 1 & 1 \\ S_0^+ & S_0^- \end{pmatrix}.$$

Then we find that

$$\begin{pmatrix} 0 & 1 \\ -D^{-1}(Q + \rho G \hat{\tau}^2) & -D^{-1}(R + R^T) \end{pmatrix} = P^{-1} \begin{pmatrix} S_0^+ & 0 \\ 0 & S_0^- \end{pmatrix} P + \text{l.o.t.}$$

It follows that

$$\begin{pmatrix} v_+ \\ v_- \end{pmatrix} = P \begin{pmatrix} v \\ hD_s v \end{pmatrix}$$

satisfies

$$(hD_s - S_0^+)v_+ = 0 \mod \mathcal{O}(h\mathcal{S}(1))$$

and

$$(hD_s - S_0^-)v_- = 0 \mod \mathcal{O}(h\mathcal{S}(1)).$$

We can view $u_+ = \mathcal{L}^{-1}v_+$ ($u_- = \mathcal{L}^{-1}v_-$) as representing incoming (outgoing) waves. This identification is justified by noticing that, for example, $v_+(\cdot, \tau) = (\mathcal{L}u_+)(\cdot, \tau)$ is exponentially decaying with increasing s for $\Re \tau > 0$.

We emphasize that the transformed DN map Λ , here, is different from the one introduced before, that is,

$$\Lambda^\tau = h\mathcal{L}\Lambda_T\mathcal{L}^{-1}.$$

With the relation between the normal directive D_s and the DN map Λ ,

$$hD_s = iD^{-1}(\Lambda^\tau - R^T) \mod \mathcal{O}(h\mathcal{S}(1)),$$

we have

$$\begin{pmatrix} v_+ \\ v_- \end{pmatrix} = P \begin{pmatrix} v \\ iD^{-1}(\Lambda^\tau - R^T)v \end{pmatrix} \mod \mathcal{O}(h\mathcal{S}(1))$$

on the boundary. We note that

$$P = (S_0^- - S_0^+)^{-1} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} P^{-\star} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$

with

$$P^{-\star} = \begin{pmatrix} 1 & S_0^- \\ 1 & S_0^+ \end{pmatrix},$$

where $P^{-\star} = P^{-\star}$ for real τ .

After we have identified the elastic parameters on the boundary, we will then know R in (3.3). Then we can have a decomposition into incoming and outgoing wave constituents (u^+ and u^- , respectively) on the boundary using Λ_T and R . Seismic imaging (inverse scattering), array receiver functions, and tomography (also using free-surface multiple scattering) all rely on this decomposition. Discussion of seismic migration and inversion schemes based on this decomposition can be found in [6, 9].

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